

Observed Decline in Planetary Albedo and Recent Global Warming Acceleration

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Key Points:

- Planetary albedo derived from CERES observations declined after 2010 $\sim 0.37\%$ per decade, increasing absorbed shortwave radiation.
- The cumulative 2000–2024 top-of-atmosphere energy imbalance reached approximately 180 ZJ and increased more rapidly after 2014
- A simple energy-balance shows that the recent acceleration of global mean surface temperature tracks observed post-2010 albedo decline

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Abstract

Recent satellite observations indicate that Earth’s planetary albedo has declined since 2010, increasing solar-energy absorption by the climate system. Using CERES top-of-atmosphere radiative flux measurements, we estimate a mean albedo decline of $\sim 0.37\%$ per decade, corresponding to an additional absorbed shortwave flux approaching 1 W m^{-2} . The associated cumulative heat uptake accelerated after 2014 and reached $\sim 180 \text{ ZJ}$ by 2024. We examine whether the observed albedo decline is dynamically consistent with the recent acceleration of global warming. A physically constrained energy-balance model combining greenhouse-gas radiative forcing with CERES-observed shortwave anomalies reproduces observed GMST evolution between 1880 and 2025 more closely than a greenhouse-gas-only baseline. The inferred attenuation timescale suggests substantial upper-ocean heat uptake during the CERES era. The observed post-2010 albedo decline appears energetically consistent with the post-2014 acceleration of GMST.

Plain Language Summary

Earth reflects part of the incoming sunlight back to space. This reflectivity, called planetary albedo, has declined since about 2010, allowing the Earth system to absorb more solar energy. Satellite observations from the CERES mission show that this decline accelerated after 2014 and was accompanied by rapidly increasing ocean heat uptake and rising global temperatures. Planetary albedo is not a fixed property of Earth. It changes together with the climate system through shifts in clouds, sea ice, water vapor, atmospheric circulation, aerosols, and ocean-atmosphere interactions such as El Niño. Because these processes are interconnected, some statistical methods used to remove short-term climate variability may also remove part of the underlying warming signal. We show that including satellite-observed albedo changes substantially improves agreement between modeled and observed global mean surface temperature from 1880 to 2025 compared with greenhouse-gas forcing alone. The results suggest that declining planetary albedo has become an important contributor to recent global warming.

1 Introduction

Recent years have seen a surge of record-breaking global temperatures and increasing evidence that global warming may be accelerating (Foster & Rahmstorf, 2011, 2026). Concurrently, anomalously warm subsurface waters have developed across the equatorial Pacific in 2026 (Figure S1), suggesting enhanced upper-ocean heat accumulation. This large-scale oceanic anomaly is used here only as qualitative supporting evidence for elevated ocean heat content, not as a basis for the statistical attribution developed below.

Further evidence for recent warming acceleration was presented by (Foster & Rahmstorf, 2026, FR26), who statistically separated the global surface air temperature signal from the effects of El Niño variability, volcanic aerosols, and solar forcing. However, their inference remains conditional on the statistical assumptions underlying the deconvolution and on the separation of internally generated variability from externally forced warming.

Mixed-baseline forcing and the albedo offset problem. The standard representation, e.g., Arias et al. (2021), of global-mean surface air temperature (GMST) change is a linear superposition of three time-dependent radiative forcing components,

$$\Delta T(t) = \lambda \left[5.35 \ln \left(\frac{c(t)}{c_{\text{PI}}} \right) + \Delta F_{\text{SW},\downarrow}(t) + \Delta F_{\text{other}}(t) \right], \quad (1)$$

where c_{PI} is the preindustrial CO_2 concentration (typically 1850), $\Delta F_{\text{SW},\downarrow}(t)$ denotes shortwave (albedo-related) solar forcing (see Eq. (S.1)), and $\Delta F_{\text{other}}(t)$ aggregates all re-

maintaining forcings (longwave greenhouse gases, aerosols, etc.). The IPCC AR6 assessment (WGI, Ch. 4 and 7) implies the climate feedback parameter

$$\lambda \approx 0.5\text{--}1.0 \text{ K (W m}^{-2}\text{)}^{-1}, \quad (2)$$

corresponding to the equilibrium climate sensitivity $ECS \approx 2.0\text{--}4.5 \text{ }^\circ\text{C}$.

In practice, however, $F_{\text{SW},\downarrow}(t)$ is not referenced to the same baseline as the logarithmic CO_2 forcing. Observational estimates of planetary albedo and absorbed solar radiation are derived primarily from the satellite era (e.g., CERES), with anomalies typically defined relative to $\sim 2000\text{--}2010$. This mismatch of reference states introduces an implicit inconsistency in Eq. (1).

Recent studies have shown that increased absorbed shortwave radiation contributed substantially to the recent rise in Earth’s energy imbalance and to the record warmth of 2023 (Goessling et al., 2025, Loeb et al., 2021).

A consistent formulation requires explicit recognition of the mixed baselines. Writing the temperature response relative to a common preindustrial reference while retaining observational albedo anomalies yields

$$\Delta T(t) = \lambda \left[5.35 \ln \left(\frac{c(t)}{c_{1850}} \right) - \frac{S_{\oplus}}{4} (\alpha(t) - \bar{\alpha}) + F_{\text{offset}} \right], \quad (3)$$

where S_{\oplus} is the solar constant, α is planetary albedo, and F_{offset} is a constant term that absorbs the unknown shortwave forcing between 1850, and the albedo reference period for $\bar{\alpha}$ is ($\sim 2000\text{--}2010$).

Without F_{offset} , Eq. (3) implicitly assumes that $\bar{\alpha}_{1850} = \bar{\alpha}_{2010}$, which is unjustified given known changes in aerosols, land use, cryosphere, and cloud fields over the industrial period. Hansen et al. (2023) acknowledge the role of declining albedo in global warming, but treat albedo change as either forcing (aerosols) or feedback (clouds), and avoid anchoring planetary albedo to preindustrial conditions.

Pre-satellite shortwave forcing components of F_{offset} . The planetary albedo term is observationally well constrained in the satellite era and contributes substantially to recent forcing trends. However, its historical baseline is not directly observed, and thus its integration into a preindustrial-referenced forcing framework necessarily introduces an unresolved offset. This mixed-baseline problem is generally avoided in the standard literature but becomes explicit in formulations such as Eq. (3).

The best-constrained component of historical shortwave (albedo-related) forcing is anthropogenic aerosol forcing. According to Arias et al. (2021) (WGI, Ch. 7), the effective radiative forcing (ERF) from aerosols over the industrial period (1750–2019) is

$$F_{\text{aer}} \approx -1.1 \text{ W m}^{-2}, \quad \text{likely range: } [-1.7, -0.4] \text{ W m}^{-2}. \quad (4)$$

This forcing is dominated by shortwave processes – direct scattering and, more importantly, aerosol–cloud interactions that increase planetary reflectivity (cloud brightening). In the present framework, F_{aer} constitutes a negative contribution to the offset term F_{offset} , partially compensating greenhouse-gas forcing prior to the satellite era.

Land-use change (LUC) also contributes to planetary albedo through surface reflectivity changes. The IPCC AR6 assessment gives

$$F_{\text{LUC}} \approx -0.15 \text{ W m}^{-2}, \quad (5)$$

for the net brightening effect of croplands and managed landscapes replacing darker natural vegetation. Although small in magnitude, this contribution is directionally consistent with aerosol forcing and thus reinforces the negative component of F_{offset} .

In contrast, cryospheric changes (snow and sea-ice albedo) introduce a positive short-wave forcing:

$$F_{\text{cryo}} \sim +0.1 \text{ to } +0.3 \text{ W m}^{-2} \text{ by } \sim 2000. \quad (6)$$

83 However, this term is poorly constrained prior to the satellite era (pre-1979), with re-
84 constructions relying on sparse observations and model inference.

85 *Implications.* If the observed CERES-era decline in planetary albedo were ex-
86 trapolated backward over the full industrial period, it would imply a positive shortwave
87 offset that would need to be reconciled with the negative aerosol forcing assessed in AR6.
88 This tension suggests that much of the observed albedo decline may be a relatively re-
89 cent phenomenon.

90 Explicit introduction of F_{offset} exposes a lack of closure in the historical energy bud-
91 get and forces reconciliation between model-derived effective radiative forcing (e.g., aerosols
92 and land use) and observation-based constraints from CERES-era radiative fluxes.

93 2 Methods

94 All data sets used in this paper, together with the software employed to generate
95 the figures, are described in the Data Availability Statement. We use the NASA CERES
96 data set (Levitus et al., 2017, Loeb et al., 2018) to compute the top-of-atmosphere (TOA)
97 planetary albedo, the Berkeley Earth data for the Global Mean Surface Temperature (GMST)
98 anomaly, and the mixed ice-core and NOAA data for the CO₂ concentration records back
99 to 1850. For at least 45 years, the CO₂ concentration in the atmosphere has been a quadratic
100 function of time (Patzek, 2026b). Although CERES EBAF products are among the most
101 carefully calibrated TOA flux data sets available, residual instrumental and intercalibra-
102 tion uncertainties remain, especially when estimating long-term trends from a relatively
103 short satellite record. The three TOA radiative flux components and the raw albedo are
104 shown in Figures S2 – S4.

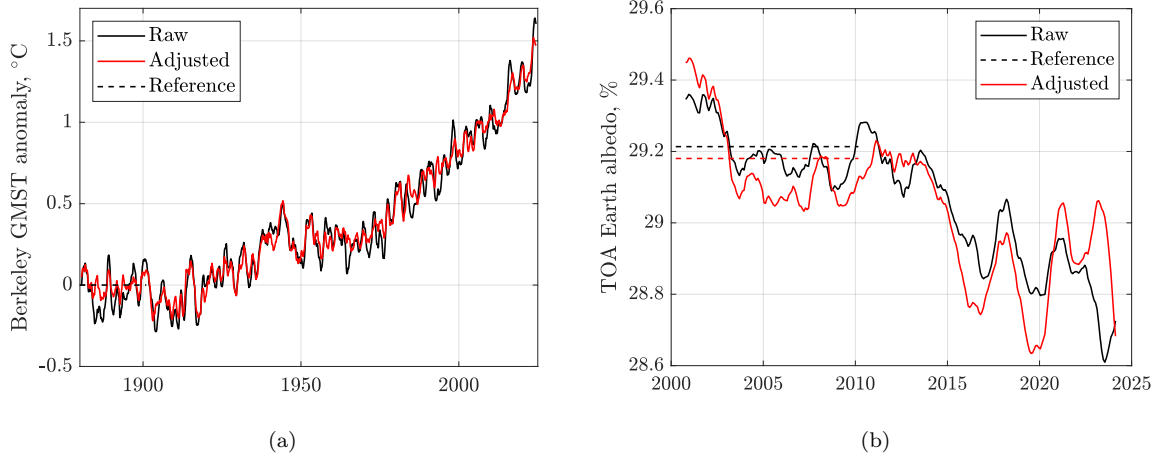


Figure 1: (a) Raw and adjusted Berkeley Earth Global Mean Surface Temperature (GMST) anomalies relative to the 1880-1900 average temperature of $-0.3 \text{ }^\circ\text{C}$. Note the thus shifted GMST anomaly was just shy of $+1.5 \text{ }^\circ\text{C}$ in 2024 and 2025. (b) The top of atmosphere (TOA) Earth albedo. The March 2000–Feb 2010 average albedo ($\approx 29.2\%$) is the broken line.

105 A 12-month running mean of the albedo anomaly, defined relative to the 2000–2010
 106 mean, and the corresponding absorbed shortwave flux are presented in Figure S5(top).
 107 The albedo-driven global cumulative heat uptake (CHU) is shown in Figure S5(bottom),
 108 together with ocean heat content (OHC) referenced to January 2000 to enable direct com-
 109 parison of trends.

110 Following the pronounced post–March 2014 increase in slope, coincident with the
 111 acceleration of global warming reported by [FR26], the CHU curve exhibits a smaller mag-
 112 nitude (by ~ 55 ZJ as of July 2024) but a steeper trend than the OHC curve. This di-
 113 vergence indicates that albedo-decline-driven CHU is becoming an increasingly impor-
 114 tant contributor to the ongoing accumulation of heat in the climate system.

115 The auxiliary data used to remove the effects of El Niño (ENSO3.4), volcanic aerosols,
 116 and solar variability from the CERES albedo signal are taken from the `ObservedData.csv`
 117 file distributed with [FR26]. The original R code was translated into MATLAB R24b
 118 and adapted for TOA albedo analysis because shortwave reflection is strongly modulated
 119 by ENSO through changes in marine stratocumulus, tropical convection, water vapor,
 120 and cloud cover (Allan & Liu, 2014, Trenberth & Caron, 2001). In addition, the inferred
 121 trend depends sensitively on the chosen start date and on whether ENSO, volcanic re-
 122 covery, and solar-cycle variability are statistically removed (Deser et al., 2012, Foster &
 123 Rahmstorf, 2011). Since the CERES record is short and autocorrelated, ordinary least-
 124 squares confidence intervals may also be overly optimistic (Santer et al., 2000, Wilks, 2011).

125 At the same time, statistically adjusting the TOA albedo is physically nontrivial.
 126 A CERES-derived decline in reflected shortwave radiation represents a real change in
 127 the TOA energy budget, but it is not necessarily an external radiative forcing in the AR6
 128 sense because part of the signal may arise from feedbacks or internally generated vari-
 129 ability rather than from an imposed forcing (Gregory et al., 2004). Moreover, the CERES
 130 record spans only 2000–2024, far shorter than the temperature records analyzed by [FR26],
 131 making robust separation of secular trends from internal variability difficult. The sta-
 132 tistical relationship between ENSO and planetary albedo therefore remains uncertain,
 133 and regression-based adjustments may inadvertently remove physically meaningful co-
 134 variance or introduce spurious trends.

Following [FR26], we adopt an additive decomposition in which the albedo trend
 is represented by a smooth function of time, $g(t)$. Owing to the large heat capacity of
 the oceans and the inertia of sea-surface temperature (SST), $g(t)$ is expected to vary slowly
 and to be largely free of high-frequency fluctuations. The model is

$$x(t) = g(t) + c_V V_{k_1}(t) + c_N N_{k_2}(t) + c_S S_{k_3}(t) + \varepsilon(t). \quad (7)$$

135 Here, $V_{k_1}(t)$ denotes the volcanic aerosol optical depth (AOD) lagged by k_1 months, $N_{k_2}(t)$
 136 is the NINO3.4 index (ENSO) lagged by k_2 months, and $S_{k_3}(t)$ is the total solar irra-
 137 diance (TSI) lagged by k_3 months. The coefficients c_V , c_N , and c_S quantify the linear
 138 response to these forcings, while $\varepsilon(t)$ is a stationary residual process.

139 The coefficients and optimal lags (k_1, k_2, k_3) are estimated by ordinary least squares
 140 following [FR26].

141 3 Results

142 The raw and adjusted GMST and albedo records are plotted in Figure 1. The GMST
 143 reference temperature differs slightly from that used by [FR26], but the resulting adjust-
 144 ments are nearly identical. The lag intervals used for the GMST and albedo adjustments
 145 are 1950–2025 and 2000–2024, respectively. With this choice of lag interval, the albedo
 146 adjustment is substantial, yet it preserves both the reference albedo, $\bar{\alpha}$, in Eq. (S.6), and
 147 the total cumulative ocean heat uptake calculated from Eq. (S.8).

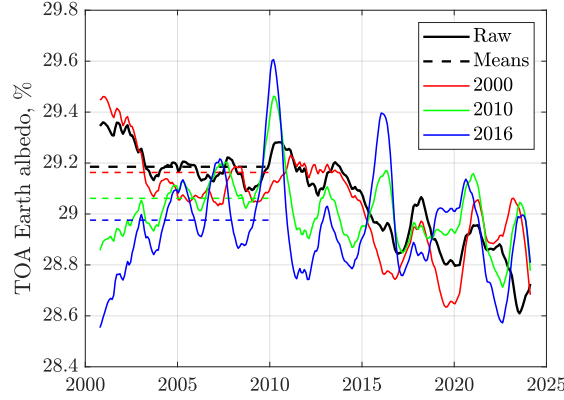


Figure 2: Twelve-month averages of the TOA albedo: the raw profile and three adjustments with the lag intervals starting in the year indicated by the legend. The reference means from 03/2000 to 2/2010 are shown as broken lines.

148 The sensitivity and instability of the albedo adjustment are illustrated in Figure 2.
 149 Only when the adjustment interval spans the full CERES record does the adjustment
 150 preserve the March 2000–February 2010 reference mean and yield a comparable cumu-
 151 lative absorbed heat anomaly. In contrast, the two alternative adjustments, initialized
 152 in February 2010 and January 2016, remain broadly in phase with one another, but dif-
 153 fer substantially from the raw albedo trajectory and produce markedly different refer-
 154 ence means.

155 Assuming that the adjusted GMST profile in Figure 1(a) provides a reasonable ap-
 156 proximation to the forced temperature response, the next question is which of the two
 157 albedo trajectories shown in Figure 1(b) yields a better fit of Equation (3) to the adjusted
 158 GMST record. To address this question, Equation (3) is recast as a constrained three-
 159 parameter optimization problem,

$$\begin{aligned}
 \min_{x \in \mathcal{F}} &= \underbrace{\sum_{i \in \mathcal{B}} [GMST(t_i) - T_{\text{base}}(x_1, x_2, t_i)]^2}_{\text{Baseline CO}_2\text{-only model}} \\
 &+ \underbrace{\sum_{j \in \mathcal{C}} [GMST(t_j) - T_{\alpha}(x_1, x_2, x_3, t_j)]^2}_{\text{Albedo change perturbation}}, \\
 T_{\text{base}} &= x_1 \left(5.35 \ln \left(\frac{c(t_i)}{c_{1850}} \right) - 1.05 + x_2 \right), \\
 T_{\alpha} &= x_1 \left[5.35 \ln \left(\frac{c(t_j)}{c_{1850}} \right) - 1.05 + x_2 + x_3 \frac{S_{\oplus}}{4} \underbrace{(\alpha(t_j) - \bar{\alpha})}_{\text{Polynomial fit}} \right], \\
 x &= [x_1, x_2, x_3] = [\lambda, F_{\text{clouds, residual}}, \alpha_{\text{mult}}].
 \end{aligned} \tag{8}$$

160 Here, \mathcal{B} denotes the set of monthly GMST observations between 1880 and 2025, while
 161 \mathcal{C} denotes the subset of monthly records for which both GMST and CERES albedo data
 162 are available (2000–2024). The feasible set \mathcal{F} imposes inequality constraints on the pa-
 163 rameter vector x , consistent with Equation (2),

$$\begin{aligned} \lambda &\in [0.5, 1] \text{ K W}^{-1} \text{ m}^2, \\ F_{\text{clouds, residual}} &\in [0.5, 1.5] \text{ W m}^{-2}, \\ \alpha_{\text{mult}} &\in [0.2, 1]. \end{aligned} \tag{9}$$

164 The constraint in Equation (9)₃ requires further explanation. It reflects the assumption
 165 that, over the relatively short CERES observation interval, the full magnitude of the TOA
 166 radiative imbalance cannot be instantaneously expressed in surface temperature because
 167 of the large heat capacity of the oceans and vigorous turbulent mixing within the up-
 168 per ocean. Consequently, the surface-temperature response to albedo-driven forcing is
 169 expected to remain approximately in phase with the TOA signal, but attenuated in am-
 170 plitude. The attenuation factor is represented here by a constant multiplier, α_{mult} , whose
 171 lower bound of 0.2 is necessarily heuristic.

172 3.1 Statistical procedure

173 The CO₂-only model represents the smooth greenhouse-gas-dominated warming
 174 trajectory using atmospheric CO₂ concentration as a compact proxy for the long-term
 175 anthropogenic forcing. It is not intended to reproduce the complete AR6 forcing decom-
 176 position, but rather to provide a deliberately parsimonious baseline against which the
 177 incremental explanatory power of CERES-observed shortwave anomaly can be tested.
 178 The combined aerosol and land-use-change (LUC) forcing is represented by a fixed off-
 179 set of -1.05 W m^{-2} , while the remaining unresolved shortwave contribution is captured
 180 by the optimized parameter $F_{\text{clouds, residual}}$. The corresponding optimal fit is shown in-
 181 Figure 3(top).

182 The albedo-augmented model additionally includes the effect of planetary albedo
 183 decline. Because the green baseline CO₂-only model temperature fit in Figure 3(top) is
 184 smooth, the raw CERES albedo anomaly in Figure 1(b) must also be replaced by a smooth
 185 polynomial approximation. Otherwise, high-frequency variability in the raw albedo record
 186 would bias the correlation analysis relative to the smooth CO₂-only fit.

187 This procedure is illustrated in Figure S6, where the raw CERES albedo anomaly
 188 is replaced by a smooth fitted curve (blue), making it statistically comparable to the smooth
 189 baseline CO₂-only model-fit (green). Both modeled temperature trajectories are then
 190 correlated against the reference GMST response (red), and the question becomes which
 191 model provides the statistically superior fit at the 95% confidence level.

192 The three time series are cross-correlated using Pearson linear correlation coeffi-
 193 cients, r_{ij} , where $i, j = \text{GMST, CO}_2, \text{Albedo}$ and $i \neq j$. The dependent correlations
 194 are then compared using the Williams test Williams (1959). Because both GMST and
 195 CERES albedo records exhibit substantial temporal autocorrelation, effective sample sizes
 196 are smaller than the nominal monthly record length. Consequently, formal significance
 197 estimates should be interpreted cautiously.

198 The albedo-augmented model, which incorporates the observed decline in plane-
 199 tary albedo, exhibits a stronger correlation with GMST than the baseline CO₂-only model
 200 (Table S1).

201 However, repeating the same procedure with the statistically adjusted albedo anomaly
 202 shown in Figure 1(b) leads to rejection of the albedo-augmented models. The adjusted
 203 albedo fit correlates with GMST less strongly than the raw albedo fit. This result sug-
 204 gests that part of the variability removed during the statistical adjustment may repre-
 205 sent physically meaningful covariance between planetary albedo and global warming, rather
 206 than purely exogenous noise.

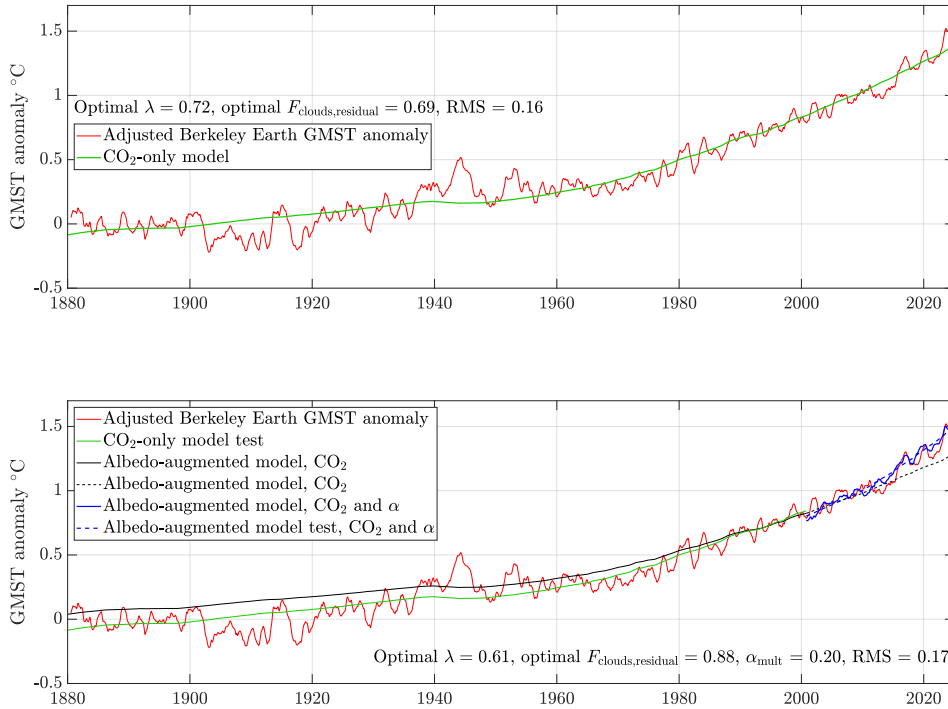


Figure 3: **(Top)** The CO₂-only model assumes that the GMST response is driven solely by the increasing atmospheric concentration of CO₂. **(Bottom)** The albedo-augmented model assumes that the GMST response is driven by the increasing atmospheric concentration of CO₂ and the decreasing raw albedo.

207 Figure 3(bottom) shows the optimal fit of the albedo-augmented model. The 12-
 208 month moving average of the raw CERES albedo anomaly (blue) is plotted together with
 209 its quadratic trend fit (broken blue curve). Interestingly, the CO₂-only temperature res-
 210 sponses between 1880 and 2000 are statistically indistinguishable. Over this interval, the
 211 baseline CO₂-only model trajectory (green) lies slightly below the albedo-augmented model
 212 trajectory (black), but the corresponding Student’s statistic satisfies $t \approx 0$. Thus, both
 213 temperature trends are effectively collinear with the observed GMST response and can-
 214 not be distinguished statistically over the 120 years of pre-CERES period. During 1880–
 215 2000, the root-mean-square (RMS) deviations of both hypotheses differ from GMST by
 216 $\mathcal{O}(10^{-1})$ °C.

217 4 Discussion

218 After 2000, the albedo-augmented model fits observed GMST significantly better
 219 than the greenhouse-gas proxy model alone. This is consistent with the interpretation
 220 that the recent warming acceleration may not be fully captured by increased longwave
 221 greenhouse trapping alone, but is also associated with enhanced absorption of incom-
 222 ing solar radiation as planetary albedo declines.

223 Since the adjusted GMST correlates more strongly with the raw albedo decline than
 224 with the adjusted albedo, several interpretations are possible: (i) the adjustment pro-
 225 cedure removes physically meaningful covariance; (ii) ENSO partly mediates the albedo

Table 1: Optimal parameters for the baseline CO₂-only and albedo-augmented models of GMST response 1880–2025.

Hypothesis	Climate parameter λ	$F_{\text{clouds,residual}}$	α_{mult}
Baseline CO ₂ -only	0.72	0.69	n/a
Albedo-augmented	0.61	0.88	0.20

226 response rather than merely contaminating it; (iii) the Foster–Rahmstorf-style decom-
 227 position is overly simplistic for TOA albedo dynamics; or (iv) the raw albedo is strongly
 228 coupled to upper-ocean temperature, remains approximately in phase with the adjusted
 229 GMST, and its surface manifestation is damped to only $\sim 20\%$ of the TOA amplitude
 230 by oceanic low-pass filtering, reducing the physical interpretability of an ENSO adjust-
 231 ment. Among these possibilities, interpretation (iv) appears broadly consistent with the
 232 observed phase coherence between GMST and the raw CERES albedo anomaly.

233 A convenient representation of the climate system response to a changing albedo
 234 is the one-box energy-balance low-pass filter

$$C \frac{dT}{dt} + \frac{1}{\lambda} T = \Delta F_{\alpha}(t), \quad (10)$$

where C is the effective heat capacity of the upper ocean–climate system, λ is the cli-
 mate feedback parameter, and $\Delta F_{\alpha}(t)$ is the time-dependent radiative forcing anomaly
 from albedo change. For an effective ocean depth of $H = 300$ m,

$$C = \rho_w c_p H \approx (1025)(3990)(300) \approx 1.23 \times 10^9 \text{ J m}^{-2} \text{ K}^{-1}.$$

The impulse response of this system is exponential, with characteristic time con-
 stant

$$\tau = \lambda C. \quad (11)$$

235 For monthly forcing data $\Delta F_{\alpha}(t)$ in W m^{-2} , a discrete-time recursion was imple-
 236 mented in MATLAB. Assuming $1/\lambda = 1.2 \text{ W m}^{-2} \text{ K}^{-1}$ (our best-fit value), the corre-
 237 sponding response times are approximately

$$\begin{aligned} H = 100 \text{ m} &\Rightarrow \tau \approx 10 \text{ yr}, \\ H = 300 \text{ m} &\Rightarrow \tau \approx 30 \text{ yr}, \end{aligned} \quad (12)$$

238 Increasing H mimics stronger coupling with the upper ocean and damps ENSO-scale vari-
 239 ability more effectively.

240 Figure 4 summarizes the central result of this paper. The left panel shows the GMST
 241 anomaly driven by albedo-change forcing after attenuation by near-surface ocean water
 242 layers of different depths, ranging from 50 to 1000 m (Equation (10)). Also shown is the
 243 TOA raw albedo anomaly calculated from the second term in Equation (3) and scaled
 244 by α_{mult} from Table 1. The results suggest that the attenuation of the TOA forcing am-
 245 plitude can be explained by heat transfer into the well-mixed upper ocean to a depth of
 246 up to 300 m, fully consistent with the timespan of the CERES observations.

247 The right panel shows the corresponding cumulative ocean heat uptake associated
 248 with the filtered CERES-observed shortwave anomaly for different ocean depths, together

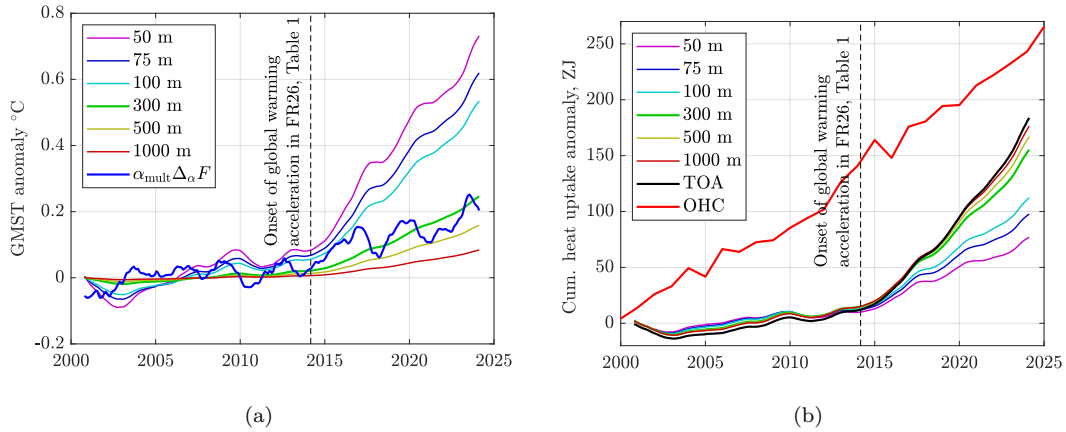


Figure 4: (a) Modeled GMST response to observed planetary albedo decline for different effective upper-ocean depths. (b) Corresponding cumulative heat uptake anomalies compared with observed ocean heat content (OHC).

249 with the TOA cumulative heat uptake from Eq. (S.7). Over the 24-year CERES record,
 250 an ocean layer approximately 300 m deep absorbs most of the TOA $\Delta F_{\alpha}(t)$ forcing.

251 Notice the pronounced increase in the slope of cumulative heat uptake after 2014,
 252 emblematic of the acceleration of global warming associated with the post-2010 decline
 253 in planetary albedo. This result is consistent with the independent statistical analysis
 254 of GMST by [FR26], which identified a similar acceleration without invoking the CERES-
 255 observed shortwave anomaly. The present results are physically consistent with the statisti-
 256 cal conclusions of [FR26].

257 Several limitations should be emphasized. First, the CERES record is short relative
 258 to multidecadal internal variability, so trend estimates remain sensitive to the chosen
 259 start date and statistical adjustment. Second, the albedo term combines forcing, feed-
 260 back, and internally generated variability and therefore is not an external forcing in the
 261 strict AR6 effective-radiative-forcing sense. Third, the one-box upper-ocean filter is a
 262 deliberately reduced representation of ocean heat uptake. The result should therefore
 263 be interpreted as evidence that declining albedo is dynamically and energetically consis-
 264 tent with recent warming acceleration, not as a complete attribution of that accelera-
 265 tion.

266 Although planetary albedo is not an external forcing independent of climate, the
 267 observed post-2010 decline appears to contribute to and amplify ongoing global warm-
 268 ing, and is energetically consistent with the post-2014 acceleration of GMST.

269 Conflict of Interest

270 The author declares no conflicts of interest relevant to this study.

271 Availability Statement

272 The input data and MATLAB software necessary to fully reproduce all calculations
 273 and figures in this paper are in Patzek (2026a), <https://doi.org/10.5281/zenodo.20432638>.

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